

GENERAL RANDERS MECHANICAL SYSTEMS

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Abstract. The general Randers spaces were introduced by R. Miron [2]. These are some generalizations of Randers spaces denoted $GR^n = (M, F + \beta)$, equipped with the Lorentz nonlinear connection. In the present paper we define the General Randers Mechanical System as a triple (M, T, F_e) , where T is the energy of the GR^n space. We obtain the expressions for the curvature and the torsion of GR^n and we give the formula for the local coefficients of the canonical connection.

1. PRELIMINARIES ON FINSLER SPACES

Let M be a n -dimensional C^∞ manifold. Denote by (TM, τ, M) the tangent bundle of M . Let $\tilde{F}^n = (M, \tilde{F}(x, y))$ be a Finsler space where $F : TM \rightarrow \mathbb{R}$ is its fundamental function and the Hessian given by $\tilde{g}_{ij}(x, y) = \frac{1}{2} \frac{\partial^2 \tilde{F}}{\partial y^i \partial y^j}$ called the fundamental tensor field of \tilde{F}^n is positive defined. The Cartan nonlinear connection \tilde{N} of the space \tilde{F}^n has the coefficients $\tilde{N}_j^i = \frac{1}{2} \frac{\partial}{\partial y^j} \left(\tilde{\gamma}_{kh}^i(x, y) y^k y^h \right)$, where we denoted by $\tilde{\gamma}_{kh}^i$ the Christoffel symbols of the metric tensor field \tilde{g}_{ij} . The Cartan nonlinear connection determines the horizontal distribution which is supplementary to the vertical distribution. The adapted basis to this distribution is $\left(\frac{\tilde{\delta}}{\delta x^i}, \frac{\partial}{\partial y^i} \right)$ with $\frac{\tilde{\delta}}{\delta x^i} = \frac{\partial}{\partial x^i} - \tilde{N}_j^i \frac{\partial}{\partial y^j}$ and the dual adapted basis is given by $(dx^i, \tilde{\delta} y^i)$, with $\tilde{\delta} y^i = dy^i + \tilde{N}_j^i dx^j$.

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The Cartan connection $C\Gamma(N) = \left(\tilde{F}_{jk}^i, \tilde{C}_{jk}^i \right)$ is given by

$$\begin{cases} \tilde{F}_{jk}^i = \frac{1}{2} \tilde{g}^{is} \left(\frac{\tilde{\delta} \tilde{g}_{sj}}{\tilde{\delta} x^k} + \frac{\tilde{\delta} \tilde{g}_{sk}}{\tilde{\delta} x^j} - \frac{\tilde{\delta} \tilde{g}_{jk}}{\tilde{\delta} x^s} \right) \\ \tilde{C}_{jk}^i = \frac{1}{2} \tilde{g}^{is} \left(\frac{\partial \tilde{g}_{sj}}{\partial y^k} + \frac{\partial \tilde{g}_{sk}}{\partial y^j} - \frac{\partial \tilde{g}_{jk}}{\partial y^s} \right). \end{cases}$$

2. GENERAL RANDERS SPACES

Let $\tilde{F}^n = \left(M, \tilde{F}(x, y) \right)$ be a Finsler space and $\beta(x, y) = b_i(x) y^i$ an 1-forms field on TM , where $b_i(x)$ is a covector field on M or on an open set of M . We shall consider the real function $L : TM \rightarrow \mathbb{R}$, $L(x, y) = \tilde{F}(x, y) + \beta(x, y)$. The pair $GR^n = (M, L(x, y))$ is called the General Randers space. The tensor field g_{ij} of GR^n is $g_{ij} = \frac{1}{2} \frac{\partial^2 L^2}{\partial y^i \partial y^j}$. Let us consider the nonlinear connection N whose local coefficients are $N_j^i = \tilde{N}_j^i - F_j^i$, where \tilde{N}_j^i is the Cartan nonlinear connection of the Finsler space \tilde{F}^n and $F_j^i = g^{ik}(x, y) F_{kj}(x)$, with $F_{kj}(x) = \frac{\partial b_j}{\partial x^k} - \frac{\partial b_k}{\partial x^j}$, the electromagnetic tensor field of the electromagnetic potentials $b_i(x)$. N is called the Lorentz nonlinear connection of the space GR^n . The local basis adapted to the Lorentz connection is $\left(\frac{\delta}{\delta x^i}, \frac{\partial}{\partial y^i} \right)$, $\frac{\delta}{\delta x^i} = \frac{\tilde{\delta}}{\delta x^i} + F_i^j \frac{\partial}{\partial y^j}$.

It is well known that using the Lorentz nonlinear connection it can be constructed an unique d-connection $D\Gamma(N) = (L_{jk}^i, C_{jk}^i)$, called the canonical metrical d-connection of the General Randers space, with the properties:

$$1. \nabla_k^H g_{ij} = 0;$$

$$2. \nabla_k^V g_{ij} = 0;$$

$$3. T_{jk}^i = 0;$$

$$4. S_{jk}^i = 0.$$

Its coefficients are given by the generalized Christoffel symbols:

$$\begin{cases} L_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\delta g_{sk}}{\delta x^j} + \frac{\delta g_{js}}{\delta x^k} - \frac{\delta g_{jk}}{\delta x^s} \right) \\ C_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\partial g_{sk}}{\partial y^j} + \frac{\partial g_{js}}{\partial y^k} - \frac{\partial g_{jk}}{\partial y^s} \right) \end{cases}$$

3. GENERAL RANDERS MECHANICAL SYSTEMS

Definition 3.1. A General Randers Mechanical System (GRMS) is a triple $\sum_{GR} = (M, T, F_e)$, where:

- i) M is an n -dimensional, real, differentiable manifold, called configuration space;
- ii) $T = y^i \frac{\partial L^2}{\partial y^i} - L^2 = L^2$ is the energy of the General Randers space GR^n ;
- iii) $F_e(x, y) = L^i(x, y) \frac{\partial}{\partial y^i}$ are the external forces given as a vertical vector field on TM and $L_i(x, y) = g_{ij} L^j(x, y)$, $i = \overline{1, n}$ are the covariant components of the F_e .

The nonlinear connection $\overset{M}{N}$ of the GRMS has the coefficients

$$N_j^i = N_j^i - \frac{1}{4} \frac{\partial F^i}{\partial y^j},$$

where N_j^i are the local coefficients of the Lorentz nonlinear connection. So,

$$\overset{M}{N}_j^i = \tilde{N}_j^i - \left(F_j^i + \frac{1}{4} \frac{\partial F^i}{\partial y^j} \right).$$

$\overset{M}{N}$ is called the Lorentz nonlinear connection of *the* GRMS and determines the horizontal distribution which is supplementary to the natural vertical distribution on TM .

A local adapted basis to these distribution is $\left(\frac{\overset{M}{\delta}}{\delta x^i}, \frac{\partial}{\partial y^i} \right)$, $i = \overline{1, n}$, where

$$\frac{\overset{M}{\delta}}{\delta x^i} = \frac{\partial}{\partial x^i} - N_j^i \frac{\partial}{\partial y^j} = \tilde{\delta} \frac{\partial}{\delta x^i} + \left(F_j^i + \frac{1}{4} \frac{\partial F^j}{\partial y^i} \right) \frac{\partial}{\partial y^j}.$$

The dual adapted basis is $\left(dx^i, \overset{M}{\delta} y^i \right)$ with

$$\overset{M}{\delta} y^i = dy^i + N_j^i dx^j = \tilde{\delta} y^i - \left(F_j^i + \frac{1}{4} \frac{\partial F^i}{\partial y^j} \right) dx^j.$$

We calculate the curvature R_{jk}^M and the torsion T_{jk}^M of the Lorentz nonlinear connection N^M of the GRMS:

$$\begin{aligned} T_{jk}^M &= \frac{\partial N_j^i}{\partial y^k} - \frac{\partial N_k^i}{\partial y^j} = \left[\frac{\partial \tilde{N}_j^i}{\partial y^k} - \frac{\partial}{\partial y^k} \left(F_j^i + \frac{1}{4} \frac{\partial F^i}{\partial y^j} \right) \right] - \\ &- \left[\frac{\partial \tilde{N}_k^i}{\partial y^j} - \frac{\partial}{\partial y^j} \left(F_k^i + \frac{1}{4} \frac{\partial F^i}{\partial y^k} \right) \right] = 0. \end{aligned}$$

$$\begin{aligned} R_{jk}^M &= \frac{\delta N_j^M}{\delta x^k} - \frac{\delta N_k^M}{\delta x^j} = \frac{\delta}{\delta x^k} \left[\tilde{N}_j^i - F_j^i - \frac{1}{4} \frac{\partial F^i}{\partial y^j} \right] - \frac{\delta}{\delta x^j} \left[\tilde{N}_k^i - F_k^i - \frac{1}{4} \frac{\partial F^i}{\partial y^k} \right] = \\ &= \frac{\tilde{\delta}}{\delta x^i} \left(\tilde{N}_j^i - F_j^i - \frac{1}{4} \frac{\partial F^i}{\partial y^j} \right) + \left(F_j^i + \frac{1}{4} \frac{\partial F^j}{\partial y^i} \right) \frac{\partial}{\partial y^j} \left(\tilde{N}_j^i - F_j^i - \frac{1}{4} \frac{\partial F^i}{\partial y^j} \right) - \\ &- \frac{\tilde{\delta}}{\delta x^j} \left(\tilde{N}_k^i - F_k^i - \frac{1}{4} \frac{\partial F^i}{\partial y^k} \right) + \left(F_k^i + \frac{1}{4} \frac{\partial F^k}{\partial y^i} \right) \frac{\partial}{\partial y^k} \left(\tilde{N}_k^i - F_k^i - \frac{1}{4} \frac{\partial F^i}{\partial y^k} \right) = \\ &= \tilde{R}_{jk}^i - \left(\frac{\tilde{\delta} F_j^i}{\delta x^i} - \frac{\tilde{\delta} F_k^i}{\delta x^j} \right) - \frac{1}{4} \left(\frac{\tilde{\delta}}{\delta x^i} \frac{\partial F^i}{\partial y^j} - \frac{\tilde{\delta}}{\delta x^j} \frac{\partial F^i}{\partial y^k} \right) + \left(F_j^i \frac{\partial \tilde{N}_j^i}{\partial y^j} - F_k^i \frac{\partial \tilde{N}_k^i}{\partial y^k} \right) - \\ &- \left(F_j^i \frac{\partial F_j^i}{\partial y^j} - F_k^i \frac{\partial F_k^i}{\partial y^k} \right) + \frac{1}{4} \left(\frac{\partial F^j}{\partial y^i} \frac{\partial \tilde{N}_j^i}{\partial y^j} - \frac{\partial F^k}{\partial y^i} \frac{\partial \tilde{N}_k^i}{\partial y^k} \right) - \frac{1}{4} \left(\frac{\partial F^j}{\partial y^i} \frac{\partial F_j^i}{\partial y^j} - \frac{\partial F^k}{\partial y^i} \frac{\partial F_k^i}{\partial y^k} \right). \end{aligned}$$

So, we can state the following theorem:

Theorem 3.1. *The torsion T_{jk}^M of the Lorentz nonlinear connection N^M of the GRMS vanishes and the curvature tensor is given by*

$$\begin{aligned} R_{jk}^M &= R_{jk}^i \sim - \left(\frac{\tilde{\delta} F_j^i}{\delta x^i} - \frac{\tilde{\delta} F_k^i}{\delta x^j} \right) - \frac{1}{4} \left(\frac{\tilde{\delta}}{\delta x^i} \frac{\partial F^i}{\partial y^j} - \frac{\tilde{\delta}}{\delta x^j} \frac{\partial F^i}{\partial y^k} \right) + \\ &+ \left(F_j^i \frac{\partial \tilde{N}_j^i}{\partial y^j} - F_k^i \frac{\partial \tilde{N}_k^i}{\partial y^k} \right) - \left(F_j^i \frac{\partial F_j^i}{\partial y^j} - F_k^i \frac{\partial F_k^i}{\partial y^k} \right) + \\ &+ \frac{1}{4} \left(\frac{\partial F^j}{\partial y^i} \frac{\partial \tilde{N}_j^i}{\partial y^j} - \frac{\partial F^k}{\partial y^i} \frac{\partial \tilde{N}_k^i}{\partial y^k} \right) - \frac{1}{4} \left(\frac{\partial F^j}{\partial y^i} \frac{\partial F_j^i}{\partial y^j} - \frac{\partial F^k}{\partial y^i} \frac{\partial F_k^i}{\partial y^k} \right) \end{aligned}$$

We fix the Lorentz nonlinear connection of GRMS and we consider a d-connection $D \Gamma \begin{pmatrix} M \\ N \end{pmatrix} = \begin{pmatrix} M & M \\ L_{jk}^i & C_{jk}^i \end{pmatrix}$ which is uniquely determined by the following axioms:

1. $\overset{M}{\nabla}_k^H g_{ij} = 0$ (D is h-metrical);
2. $\overset{M}{\nabla}_k^V g_{ij} = 0$ (D is v-metrical);
3. $T_{jk}^i = 0$ (D is h-torsion free);
4. $S_{jk}^i = 0$ (D is v-torsion free), where

$$\overset{M}{\nabla}_k^H g_{ij} = \frac{\delta g_{ij}}{\delta x^k} - L_{ik}^s g_{sj} - L_{jk}^s g_{is}, \quad \overset{M}{\nabla}_k^V g_{ij} = \frac{\partial g_{ij}}{\partial y^k} - C_{ik}^s g_{sj} - C_{jk}^s g_{is}.$$

The local coefficients of $D \overset{M}{\Gamma} \left(\overset{M}{N} \right) = \left(L_{jk}^i, C_{jk}^i \right)$ are expressed by the generalized Christoffel symbols:

$$\begin{cases} L_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\delta g_{sj}}{\delta x^k} + \frac{\delta g_{sk}}{\delta x^j} - \frac{\delta g_{jk}}{\delta x^s} \right) \\ C_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\partial g_{sj}}{\partial y^k} + \frac{\partial g_{sk}}{\partial y^j} - \frac{\partial g_{jk}}{\partial y^s} \right). \end{cases}$$

Through a direct calculation and using the results from [2] we get the explicit form of these coefficients in the following theorem:

Theorem 3.2. *The coefficients L_{jk}^i, C_{jk}^i of $D \overset{M}{\Gamma} \left(\overset{M}{N} \right) = \left(L_{jk}^i, C_{jk}^i \right)$ are given by*

$$L_{jk}^i = L_{jk}^i + A_{jk}^i, \quad C_{jk}^i = C_{jk}^i,$$

where $CT(N) = (L_{jk}^i, C_{jk}^i)$ is the canonical metrical d-connection of GR^n , with

$$L_{jk}^i = \tilde{F}_{jk}^i + \tilde{C}_{js}^i F_k^s + B_{jk}^i, \quad C_{jk}^i = \tilde{C}_{jk}^i + E_{jk}^i$$

and

$$\begin{aligned} A_{jk}^i &= \frac{1}{4} g^{is} \left(C_{skh} \frac{\partial F^h}{\partial y^j} + C_{jsh} \frac{\partial F^h}{\partial y^k} - C_{jkh} \frac{\partial F^h}{\partial y^s} \right) \\ B_{jk}^i &= \frac{1}{2} g^{ir} \left(\tilde{\nabla}_k^H g_{jr} + \tilde{\nabla}_j^H g_{rk} - \tilde{\nabla}_r^H g_{jk} \right) + F_k^s E_{sj}^i + F_j^s C_{sk}^i - g^{ir} F_r^s C_{sjk} \\ E_{jk}^i &= -\frac{F}{L} \frac{y^i}{F} b^s \tilde{C}_{sjk}^i + \frac{1}{2F^2} g^{is} \left[\left(F b_k - \beta \frac{\partial F}{\partial y^i} \right) F \frac{\partial^2 F}{\partial y^s \partial y^j} + \right. \\ &\quad \left. + \left(F b_s - \beta \frac{\partial F}{\partial y^s} \right) F \frac{\partial^2 F}{\partial y^j \partial y^k} + \left(F b_j - \beta \frac{\partial F}{\partial y^j} \right) F \frac{\partial^2 F}{\partial y^s \partial y^k} \right] \end{aligned}$$

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