

R - MECHANICAL SYSTEMS

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Abstract. In the present paper we study a particular case of Finslerian mechanical system This is defined by a triple $\sum_R = (M, F^2, F_e)$, where M is the configuration space, $F(x, y) = \alpha(x, y) + \beta(x, y)$ is a Randers metric and $F_e = a_{jk}^i(x) y^j y^k \frac{\partial}{\partial y^i}$ are the external forces.

1. INTRODUCTION AND PRELIMINARIES

Let M be an n -dimensional, real C^∞ manifold. Denote by (TM, τ, M) the tangent bundle of M and let $F^n = (M, F(x, y))$ be a Finsler space, where $F : TM \rightarrow R_+$ is it fundamental function, i.e., F verifies the following axioms:

- i) F is a differentiable function on $\widetilde{TM} = TM - \{0\}$ and it is continuous on the null section of the projection $\tau : TM \rightarrow M$;
- ii) F is positively 1- homogeneous with respect to the variables y^i ;
- iii) For every $(x, y) \in \widetilde{TM}$ the Hessian of F^2 with respect y^i is positively defined.

Consequently, the d-tensor field $g_{ij}(x, y) = \frac{1}{2} \frac{\partial^2 F^2}{\partial y^i \partial y^j}$ is positively defined. It is called the fundamental tensor or metric tensor of F^n .

For a non-zero vector $y \in T_p M$, define $h_y(u, v) = g_y(u, v) - F^{-2}(y) g_y(y, u) g_y(y, v)$, with $u, v \in T_p M$. $h = \{h_y\}$ is called the angular metric of F . The geodesics of F are characterized locally by $\frac{d^2 x^i}{dt^2} + 2G^i(x, \frac{dx}{dt}) = 0$, where $G^i = \frac{1}{4} g^{ik} \left\{ 2 \frac{\partial g_{pk}}{\partial x^q} - \frac{\partial g_{pq}}{\partial x^k} \right\} y^p y^q$.

Keywords and phrases: Randers space, mechanical system.

(2010) Mathematics Subject Classification: 53C60, 53B40.

2. RANDERS SPACES

Let $F^n = (M, F(x, y))$ be a Finsler space with the fundamental function $F(x, y) = \alpha(x, y) + \beta(x, y)$ where $\alpha(x, y) = \sqrt{a_{ij}(x)y^i y^j}$ and $\beta(x, y) = b_i(x)y^i$. $a = a_{ij}(x)dx^i dx^j$ is a pseudo-Riemannian metric on M and $b_i(x)$ is a covector field on the manifold M . We denote by " \mid " the covariant differentiation with respect to the Levi-Civita connection $\gamma = \{\gamma_{jk}^i\}$ of the associated Riemannian space (M, α^2) . We shall use the notation as follows [9]:

$$\begin{aligned}\tilde{l}_i &= \frac{\partial \alpha}{\partial y^i} = \frac{a_{ij}y^j}{\alpha}, l_i = \frac{\partial F}{\partial y^i} = \tilde{l}_i + b_i, l^i = \frac{y^i}{F}, \tilde{l}^i = \frac{y^i}{\alpha}, \\ b_{j|k} &= \frac{\partial b_j}{\partial x^k} - b_s \gamma_{jk}^s, r_{ij} = \frac{1}{2}(b_{i|j} + b_{j|i}), s_{ij} = \frac{1}{2}(b_{i|j} - b_{j|i}), \\ s_j^i &= a^{ih} s_{hj}, s_j = b_i s_j^i, e_{ij} = r_{ij} + b_i s_j + b_j s_i, \\ e_{00} &= e_{ij} y^i y^j, s_0 = s_i y^i, s_0^i = s_j^i y^j.\end{aligned}$$

The fundamental tensor g_{ij} of the Randers metric $F = \alpha + \beta$ is expressed by

$$(2.1) \quad g_{ij} = \frac{F}{\alpha}(a_{ij} - \tilde{l}_i \tilde{l}_j) + l_i l_j.$$

The functions $G^i(x, y) = \frac{1}{2}\Gamma_{jk}^i(x, y)y^j y^k$ are the components of the geodesic spray $S = y^i \frac{\partial}{\partial x^i} - 2G^i \frac{\partial}{\partial y^i}$ and $\Gamma_{jk}^i(x, y)$ are the Christoffel symbols of the metric tensor g_{ij} . By a direct calculation

$$(2.2) \quad G^i = (\gamma_{jk}^i + l^i b_{j|k})y^j y^k + (a^{ij} - l^i b^j)(b_{j|k} - b_{k|j})\alpha y^k,$$

or

$$(2.3) \quad G^i = \tilde{G}^i + \frac{1}{2F}(r_{kl}y^k y^l - 2\alpha b_r a^{rp} s_{pl}y^l)y^i + \alpha a^{ir} s_{rl}y^l,$$

or, equivalently,

$$(2.4) \quad G^i = \tilde{G}^i + \frac{e_{00}}{2F}y^i - s_0 y^i + \alpha s_0^i,$$

with \tilde{G}^i the components of the geodesic spray of the Riemannian space.

The Cartan nonlinear connection $\overset{C}{N}$ for the Finsler space $F^n = (M, F = \alpha + \beta)$ has the coefficients

$$(2.5) \quad \overset{C}{N}_j^i = \frac{\partial G^i}{\partial y^j}.$$

Definition 2.1. The Finsler space $F^n = (M, F = \alpha + \beta)$ equipped with Cartan nonlinear connection $\overset{C}{N}$ is called a *Randers space* and it is denoted by $RF^n = \left(M, \alpha + \beta, \overset{C}{N}\right)$.

The local basis adapted to the Cartan nonlinear connection is $\left(\frac{\delta}{\delta x^i}, \frac{\partial}{\partial y^i}\right)_{i=\overline{1,n}}$ with

$$(2.6) \quad \frac{\overset{C}{\delta}}{\delta x^i} = \frac{\partial}{\partial x^i} - N_i^k \frac{\partial}{\partial y^k}.$$

Theorem 2.1. *There exists an unique $\overset{C}{N}$ -metrical connection $C\Gamma\left(\overset{C}{N}\right) = (F_{jk}^i, C_{jk}^i)$ of the Randers space RF^n which verifies the following axioms:*

- i) $\overset{C}{\nabla}_k^H g_{ij} = 0; \overset{C}{\nabla}_k^V g_{ij} = 0;$
- ii) $\overset{C}{T}_{jk}^i = 0; \overset{C}{S}_{jk}^i = 0.$

The connection $C\Gamma\left(\overset{C}{N}\right)$ has the coefficients expressed by the generalized Christoffel symbols:

$$(2.7) \quad \begin{cases} \overset{C}{F}_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\overset{C}{\delta} g_{sj}}{\delta x^k} + \frac{\overset{C}{\delta} g_{sk}}{\delta x^j} - \frac{\overset{C}{\delta} g_{jk}}{\delta x^s} \right) \\ \overset{C}{C}_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\partial g_{sj}}{\partial y^k} + \frac{\partial g_{sk}}{\partial y^j} - \frac{\partial g_{jk}}{\partial y^s} \right). \end{cases}$$

The integrability tensor of $\overset{C}{N}$ is

$$(2.8) \quad R_{jk}^i = \frac{\overset{C}{\delta} \overset{C}{N}_j^i}{\delta x^k} - \frac{\overset{C}{\delta} \overset{C}{N}_k^i}{\delta x^j}.$$

3. R-MECHANICAL SYSTEMS

For a manifold M , that is the configuration space of Finslerian dynamical system, let us consider the tangent bundle TM to which we refer to as the velocity space. Suppose that there exists a Randers metric $F = \alpha + \beta$ on \widehat{TM} and $\alpha_{jk}^i(x)$ a symmetric tensor on the configuration space M of type $(1, 2)$.

Definition 3.1. An R -mechanical system is a triple $\sum_R = (M, F^2, F_e)$, where $F_e = a_{jk}^i(x) y^j y^k \frac{\partial}{\partial y^i}$ are the external forces, and $F^n = (M, F = \alpha + \beta)$ is a Randers space endowed with Cartan non-linear connection N .

We denote $F^i(x, y) = a_{jk}^i(x) y^j y^k$ and one considers $F_i(x, y) = g_{ij} F^j(x, y)$ the covariant components of the external forces F_e .

The energy of the Randers space $RF^n = (M, F = \alpha + \beta)$ is

$$(3.1) \quad \varepsilon_{F^2} = y^i \frac{\partial(\alpha + \beta)^2}{\partial y^i} - (\alpha + \beta)^2 = (\alpha + \beta)^2 = g_{ij}(x, y) y^i y^j.$$

If the external forces F_e are global defined on \widetilde{TM} we obtain

Theorem 3.2. [8] For the Finslerian mechanical system $\sum_R = (M, (\alpha + \beta)^2, F_e)$ the following properties hold true:

i) The operator S defined by

$$(3.2) \quad S = y^i \frac{\partial}{\partial x^i} - (2G^i - \frac{1}{2} a_{jk}^i(x) y^j y^k) \frac{\partial}{\partial y^i}$$

is a vector field, global defined on the phase space \widetilde{TM} .

ii) S is a semispray which depends only on \sum_R and it is a spray if F_e are 2-homogeneous with respect to y^i .

iii) The integral curves of the vector field S are the evolution curves given by the Lagrange equations of \sum_R :

$$(3.3) \quad \frac{d^2 x^i}{dt^2} + \Gamma_{jk}^i(x, \frac{dx}{dt}) \frac{dx^j}{dt} \frac{dx^k}{dt} = \frac{1}{2} F^i(x, \frac{dx}{dt}).$$

The semispray S (3.2) has the coefficients G^i expressed by

$$(3.4) \quad 2G^i = 2G^i - \frac{1}{2} a_{jk}^i(x) y^j y^k = \Gamma_{jk}^i(x, y) y^j y^k - \frac{1}{2} a_{jk}^i(x) y^j y^k.$$

Thus, the canonical nonlinear connection N of the Finslerian mechanical system \sum_R has the coefficients

$$(3.5) \quad N_j^i = \frac{\partial G^i}{\partial y^j} = N_j^i - a_{jk}^i(x) y^k$$

The nonlinear connection N determines the horizontal distribution N which is supplementary to the natural vertical distribution on the

tangent manifold \widetilde{TM} . A local adapted basis to these distributions is $\left(\frac{\overset{R}{\delta}}{\delta x^j}, \frac{\partial}{\partial y^j}\right)_{j=\overline{1,n}}$ where

$$(3.6) \quad \begin{aligned} \frac{\overset{R}{\delta}}{\delta x^j} &= \frac{\partial}{\partial x^i} - N_j^i \frac{\partial}{\partial y^j} = \frac{\partial}{\partial x^i} - N_j^i \frac{\overset{C}{\partial}}{\partial y^j} + a_{jk}^i(x) y^k \frac{\partial}{\partial y^j} \\ &= \frac{\overset{R}{\delta}}{\delta x^i} + a_{jk}^i(x) y^k \frac{\partial}{\partial y^j}. \end{aligned}$$

The adapted cobasis is $\left(dx^i, \delta y^i\right)_{i=\overline{1,n}}$ with

$$(3.7) \quad \begin{aligned} \overset{R}{\delta} y^i &= dy^i + N_j^i dx^j = dy^i + N_j^i \overset{C}{dx}^j - a_{jk}^i(x) y^k dx^j \\ &= \delta y^i - a_{jk}^i(x) y^k dx^j. \end{aligned}$$

We determine the torsion $T_{jk}^{\overset{R}{i}}$ and the curvature $R_{jk}^{\overset{R}{i}}$ of the canonical connection $\overset{R}{N}$ by a direct calculation:

$$(3.8) \quad T_{jk}^{\overset{R}{i}} = \frac{\partial \overset{R}{N}_j^i}{\partial y^k} - \frac{\partial \overset{R}{N}_k^i}{\partial y^j} = 0.$$

$$(3.9) \quad R_{jk}^{\overset{R}{i}} = \frac{\overset{R}{\delta} \overset{R}{N}_j^i}{\delta x^k} - \frac{\overset{R}{\delta} \overset{R}{N}_k^i}{\delta x^j} = \overset{C}{R}_{jk}^i + (N_k^s a_{jk}^i - N_j^s a_{sk}^i).$$

Applying the theory from the book [9] the following theorem holds:

Theorem 3.3. *Let $\sum_R = (M, (\alpha + \beta)^2, F_e)$ be a R - mechanical system and $\overset{R}{N}$ the canonical nonlinear connection of \sum_R . There exists an unique d -connection $R\Gamma(\overset{R}{N}) = (\overset{R}{F}_{jk}^i, \overset{R}{C}_{jk}^i)$ determined by the following axioms:*

- i) $\overset{R}{\nabla}_k^H g_{ij} = 0; \overset{R}{\nabla}_k^V g_{ij} = 0;$
- ii) $\overset{R}{T}_{jk}^i = 0; \overset{R}{S}_{jk}^i = 0,$

where

$$(3.10) \quad \begin{aligned} \nabla_k^H g_{ij} &= \frac{\delta g_{ij}}{\delta x^k} - F_{ik}^s g_{sj} - F_{jk}^s g_{is} \\ \nabla_k^V g_{ij} &= \frac{\partial g_{ij}}{\partial y^k} - C_{ik}^s g_{sj} - C_{jk}^s g_{is}. \end{aligned}$$

We call this connection the canonical metrical d -connection of \sum_R .

Theorem 3.4. *The local coefficients of the canonical metrical d -connection of \sum_R are*

$$(3.11) \quad \begin{cases} F_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\delta g_{sj}}{\delta x^k} + \frac{\delta g_{sk}}{\delta x^j} - \frac{\delta g_{jk}}{\delta x^s} \right) \\ C_{jk}^i = \frac{1}{2} g^{is} \left(\frac{\partial g_{sj}}{\partial y^k} + \frac{\partial g_{sk}}{\partial y^j} - \frac{\partial g_{jk}}{\partial y^s} \right) \end{cases}$$

In order to calculate F_{jk}^i and C_{jk}^i we have from (3.6):

$$(3.12) \quad \frac{\delta g_{sj}}{\delta x^k} = \frac{C}{\delta x^k} g_{sj} + a_{jk}^i(x) y^k \frac{\partial g_{sj}}{\partial y^h}.$$

Now, the developed expression of the coefficients F_{jk}^i and C_{jk}^i is given in the next theorem:

Theorem 3.5. *The canonical metrical d -connection of \sum_R has the coefficients*

$$\begin{cases} F_{jk}^i = F_{jk}^i + \frac{1}{2} g^{is} \left(a_{kp}^h y^p \frac{\partial g_{sj}}{\partial y^h} + a_{jp}^h y^p \frac{\partial g_{sk}}{\partial y^h} - a_{sp}^h y^p \frac{\partial g_{jk}}{\partial y^h} \right) \\ C_{jk}^i = C_{jk}^i. \end{cases}$$

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