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RHEONOMIC GENERAL RANDERS MECHANICAL SYSTEMS

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Abstract. This paper presents a special type of rheonomic Finslerian Mechanical System called rheonomic General Randers Mechanical System, establishing the canonical semispray, the local coefficients of the Lorentz connection and the expressions for the curvature and for the torsion.

1. Preliminary. Rheonomic Finsler Spaces

Let M be a real n-dimensional differentiable manifold and (TM, π, M) the tangent bundle of M. We consider $E = TM \times R$ a 2n + 1- dimensional real manifold. In a local chart $U \times (a, b)$ the points $u = (x, y, t) \in E$ have the local coordinates (x^i, y^i, t) . Let $rF^n = (M, F(x, y, t))$ be a rheonomic Finsler space where $F: E \to R$ is the fundamental function and the Hessian of F given by $g_{ij}(x, y, t) = \frac{1}{2} \frac{\partial^2 F}{\partial y^i \partial y^j}$, called the fundamental tensor of rF^n , is positive defined.

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The Cartan nonlinear connection N has the coefficients $\left(N_{i}^{i}\left(x,y,t\right),N_{i}^{0}\left(x,y,t\right)\right)$ with

(1.1)
$$N_{j}^{i}(x,y,t) = \frac{1}{2} \frac{\partial}{\partial y^{j}} \left(\gamma_{kh}^{i}(x,y,t) y^{k} y^{h} \right)$$
$$N_{j}^{0}(x,y,t) = \frac{1}{2} \frac{\partial g_{jk}}{\partial t} y^{k}$$

and γ_{kj}^i are the Christoffel symbols of the fundamental tensor $g_{ij}(x, y, t)$. N determines the horizontal distribution on E which is supplementary to the vertical distribution. The adapted basis to these distribution is $\left(\frac{\delta}{\delta x^i}, \frac{\partial}{\partial y^i}, \frac{\partial}{\partial t}\right)$ with

(1.2)
$$\frac{\delta}{\delta x^{i}} = \frac{\partial}{\partial x^{i}} - N_{i}^{j}(x, y, t) \frac{\partial}{\partial y^{j}} - N_{i}^{0}(x, y, t) \frac{\partial}{\partial t}.$$

The dual adapted basis is $(dx^i, \delta y^i, dt)$ where

(1.3)
$$\delta y^{i} = dy^{i} + N_{j}^{i}(x, y, t) dx^{j} \\ \delta t = dt + N_{j}^{0}(x, y, t) dx^{j}.$$

2. Rheonomic General Randers Spaces

We consider $\beta(x, y, t) = b_i(x, t)y^i$ with $b_i(x, y)$ a covector field on $M \times R$ and the metric $L : E \to R$, $L(x, y, t) = F(x, y, t) + \beta(x, y, t)$. The pair $rGR^n = (M, L(x, y, t))$ is called a rheonomic General Randers space. The tensor field g_{ij} of rGR^n is

(2.1)
$$g_{ij}^{L}(x,y,t) = \frac{1}{2} \frac{\partial^{2} L^{2}}{\partial y^{i} \partial y^{j}}.$$

We denote

(2.2)
$$F_{ij}(x,t) = \frac{\partial b_j}{\partial x^i} - \frac{\partial b_i}{\partial x^j}$$

and we consider

(2.3)
$$F_{j}^{i}(x, y, t) = g^{ik}(x, y, t) F_{kj}(x, t).$$

Let us consider in a rGR^n space the nonlinear connection N whose local coefficients are $\begin{pmatrix} L \\ N^i_j\left(x,y,t\right), N^0_j\left(x,y,t\right) \end{pmatrix}$ with

(2.4)
$$\begin{cases} N_j^i = N_j^i - F_j^i \\ N_j^0 = N_j^0 - F_j^i \end{cases}$$

N is called Lorentz nonlinear connection of the space rGR^n . The local basis adapted to the Lorentz nonlinear connection is $\left(\frac{L}{\delta x^i}, \frac{\partial}{\partial y^i}, \frac{\partial}{\partial t}\right)$, where

(2.5)
$$\frac{\overset{L}{\delta}}{\delta x^{i}} = \frac{\delta}{\delta x^{i}} + F_{j}^{i} \frac{\partial}{\partial y^{j}} + F_{j}^{i} \frac{\partial}{\partial t}.$$

The dual adapted basis is $\left(dx^{i}, \overset{L}{\delta}y^{i}, \overset{L}{\delta}t\right)$ with

(2.6)
$$\begin{aligned}
\delta y^i &= \delta y^i - F_j^i dx^j \\
\delta t &= \delta t - F_j^i dx^j.
\end{aligned}$$

The weak torsion T_{jk}^{L} of rGR^{n} space is

(2.7)
$$T_{jk}^{L} = \frac{\partial N_{j}^{i}}{\partial y^{k}} - \frac{\partial N_{k}^{i}}{\partial y^{j}}$$

And the curvature tensor R_{jk}^{L} is

(2.8)
$$R_{jk}^{L} = \frac{\delta N_{j}^{i}}{\partial x^{k}} - \frac{\delta N_{k}^{i}}{\partial x^{j}}.$$

By a direct calculus we can state the following theorem:

Theorem 2.1. The torsion T_{jk}^i of the Lorentz nonlinear connection of the rGR^n space vanishes and the curvature tensor is given by

$$(2.9)$$

$$R_{jk}^{i} = \left(\frac{\delta N_{j}^{i}}{\delta x^{k}} - \frac{\delta N_{k}^{i}}{\delta x^{j}}\right) - \left(\frac{\delta F_{j}^{i}}{\delta x^{k}} - \frac{\delta F_{k}^{i}}{\delta x^{j}}\right) + \left(F_{j}^{i} \frac{\partial N_{j}^{i}}{\partial y^{k}} - F_{k}^{i} \frac{\partial N_{k}^{i}}{\partial y^{j}}\right) - \left(F_{j}^{i} \frac{\partial F_{j}^{i}}{\partial y^{k}} - F_{k}^{i} \frac{\partial F_{k}^{i}}{\partial y^{j}}\right) + \left(F_{j}^{i} \frac{\partial N_{j}^{i}}{\partial t} - F_{k}^{i} \frac{\partial N_{k}^{i}}{\partial t}\right) - \left(F_{j}^{i} \frac{\partial F_{j}^{i}}{\partial t} - F_{k}^{i} \frac{\partial F_{k}^{i}}{\partial t}\right).$$

3. Rheonomic General Randers Mechanical Systems

Definition 3.1. A Rheonomic General Randers Mechanical System (rhGRMS) is a triple $\sum = (M, L^2(x, y, t), \sigma_i(x, y, t))$, where:

- \mathbf{i}) M is an n-dimensional, real, differentiable manifold;
- ii)L(x, y, t) is the fundamental function of the rheonomic general CityplaceRanders space rGR^n ;
 - iii) $\sigma_i(x, y, t)$ is a d-covector field called the external force of \sum .

The evolution equations of the rhGRMS \sum are given by the system of second order differential equations:

(3.1)
$$\frac{d^{2}x^{i}}{dt} + 2 \overset{M}{G^{i}}(x, y, t) + \overset{M}{N_{0}^{i}}(x, y, t) = \frac{1}{2}\sigma^{i}(x, y, t)$$

where

(3.2)
$$\sigma^i = g^{ij}\sigma_j$$

and

(3.3)
$$G^{i} = \frac{1}{2} \gamma_{rs}^{i}(x, y, t) y^{r} y^{s}.$$

The equations (3.1) determine a semispray $\stackrel{M}{S}$, or a dynamical system on $TM \times R$.

We can prove

Theorem 3.1. a) The semispray $\stackrel{M}{S}$ on $TM \times R$, is given by

$$\overset{M}{S} = y^{i} \frac{\partial}{\partial x^{i}} - \left(2 \overset{M}{G}^{i}(x, y, t) - \frac{1}{2} \sigma^{i}(x, y, t) + N_{0}^{i} \overset{M}{(x, y, t)} \right) \frac{\partial}{\partial y^{i}} + \frac{\partial}{\partial t}.$$

- b) $\stackrel{M}{S}$ is a dynamical system on $\stackrel{\sim}{TM} \times R$ depending only of the rhGRMS \sum . We call this semispray the evolution semispray of the rhGRMS \sum .
- c) The integral curves of $\stackrel{M}{S}$ are the evolution curves of \sum given by (3.1).

The evolution semispray (3.3) determines a nonlinear connection N = N depending only by the $rhGRMS \sum$. The local coefficients are $\binom{M}{N_j}, N_j^0$, where

$$(3.4) \begin{array}{c} N_{j}^{i} = N_{j}^{i} - \frac{1}{4} \frac{\partial \sigma^{i}}{\partial y^{j}} \\ N_{j}^{i} = N_{j}^{0}, \end{array}$$

with N_j^i the local coefficients of the Lorenz nonlinear connection of the rheonomic General Randers space. So,

(3.5)
$$N_j^i = N_j^i - \left(F_j^i + \frac{1}{4} \frac{\partial \sigma^i}{\partial y^j}\right)$$
$$N_j^0 = N_j^0 - F_j^i.$$

 $\stackrel{M}{N}$ is called the Lorentz nonlinear connection of the rhGRMS \sum and determines a direct decomposition on $TM \times R$ into horizontal and vertical subspaces. The adapted basis to this decomposition is $\left(\frac{M}{\delta x^i}, \frac{\partial}{\partial y^i}, \frac{\partial}{\partial t}\right)$ where

(3.6)
$$\frac{\frac{\delta}{\delta x^i}}{\delta x^i} = \frac{\delta}{\delta x^i} + \left(F_j^i + \frac{1}{4} \frac{\partial \sigma^j}{\partial y^i}\right) \frac{\partial}{\partial y^j} + F_j^i \frac{\partial}{\partial t}.$$

The dual basis is $\left(dx^{i}, \delta^{i}, \delta^{t}\right)$ where

Theorem 3.2. The torsion T_{jk}^{M} of the Lorentz nonlinear connection N = N of the N = N vanishes and the curvature tensor N = N = N is given by

(3.8)
$$R_{jk}^{i} = \frac{\stackrel{M}{\delta} \stackrel{M}{N_{j}^{i}}}{\delta x^{k}} - \frac{\stackrel{M}{\delta} \stackrel{M}{N_{k}^{i}}}{\delta x^{j}}.$$

REFERENCES

- [1] Bucataru SnI., Miron R., Finsler-Lagrange Geometry- Applications to dynamical systems, Ed. Academiei Romane, Bucuresti, (2007).
- [2] Frigioiu C., On the geometrization of the rheonomic Finslerian Mechanical Systems, Analele Științifice ale Univ. "Al.I. Cuza", Iași, vol. LIII, (2007), 145-155.
- [3] Lungu O., Niminet V, General Randers Mechanical Systems, Scientific Studies and Research. Series Mathematics and Informatics, Vol. 22, No.1, (2012), 25-30.
- [4] Miron R., General Randers Spaces, Lagrange and Finsler Geometry, Ed. By P.L. Antonelli and R. Miron, FTPH 59, Kluwer Acad. Publ., (1996), 123-140.
- [5] Miron R., Lagrangian and Hamiltonian Geometries- Applications to Analytical Mechanics, Ed. Academiei Romane, Ed. Fair Partners, Bucuresti, (2011).
- [6] Miron R., Dynamical Systems of Lagrangian and Hamiltonian Mechanical Systems, Advance Studies in Pure Mathematics, 48, (2007), 309-340.
- [7] Miron R., Niminet V., **The Lagrangian Geometrical Model and the Associated Dynamical System of a Nonholonomic Mechanical System**, Facta Universitas, seria Mechanics, Automatic and Control robotic, vol. 5, nr. 1, 1-24.
- [8] Niminet V., New Geometrical Properties of Generalized Lagrange Lpaces of Relativistic Optics, Tensor New series, 68,1, 66-70.
- [9] Niminet V., Generalized Lagrange Spaces of Relativistics Optics, Tensor new series, 66,2, 180-184.
- [10] Niminet V., Generalized Lagrange Spaces Derived from a Finsler-Synge Metric, Libertas Mathematica, vol. 24, 33-38.

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